Automatica 155 (2023) 111124

Contents lists available at ScienceDirect

Automatica

journal homepage: www.elsevier.com/locate/automatica

On the regularization and optimization in quantum detector tomography^{*}

Shuixin Xiao ^{a,b}, Yuanlong Wang ^{c,d}, Jun Zhang ^{a,*}, Daoyi Dong ^{b,*}, Shota Yokoyama ^{b,e}, Ian R. Petersen ^f, Hidehiro Yonezawa ^{b,e}

^a University of Michigan – Shanghai Jiao Tong University Joint Institute, Shanghai Jiao Tong University, Shanghai 200240, China

^b School of Engineering and Information Technology, University of New South Wales, Canberra ACT 2600, Australia

^c Key Laboratory of Systems and Control, Academy of Mathematics and Systems Science, Chinese Academy of Sciences, Beijing 100190, China

^d Centre for Quantum Computation and Communication Technology (Australian Research Council), Centre for Quantum Dynamics, Griffith

University, Brisbane, Queensland 4111, Australia

^e Centre for Quantum Computation and Communication Technology, Australian Research Council, Canberra, ACT 2600, Australia ^f School of Engineering, Australian National University, Canberra, ACT 2601, Australia

ARTICLE INFO

Article history: Received 1 August 2022 Received in revised form 8 January 2023 Accepted 9 April 2023 Available online 17 June 2023

Keywords: Quantum system identification Quantum detector tomography Quantum system Regularization

ABSTRACT

Quantum detector tomography (QDT) is a fundamental technique for calibrating quantum devices and performing quantum engineering tasks. In this paper, we utilize regularization to improve the QDT accuracy whenever the probe states are informationally complete or informationally incomplete. In the informationally complete scenario, without regularization, we optimize the resource (probe state) distribution by converting it to a semidefinite programming problem. Then in both the informationally complete and informationally incomplete scenarios, we discuss different regularization forms and prove the mean squared error scales as O(1/N) or tends to a constant with N state copies under the static assumption. We also characterize the ideal best regularization for the identifiable parameters, accounting for both the informationally complete and informationally incomplete scenarios. Numerical examples demonstrate the effectiveness of different regularization forms and a quantum optical experiment test shows that a suitable regularization form can reach a reduced mean squared error. (2023 Elsevier Ltd. All rights reserved)

1. Introduction

In the past decades, significant progress has been achieved in a variety of fields of quantum science and technology, including quantum computation (DiVincenzo, 1995), quantum communication (Nielsen & Chuang, 2010) and quantum sensing (Degen, Reinhard, & Cappellaro, 2017). In these applications, it is often necessary to develop efficient estimation methods to acquire information about quantum systems and quantum system identification has attracted wide attention (Burgarth & Yuasa, 2012;

E-mail addresses: xiaoshuixin@sjtu.edu.cn (S. Xiao),

wangyuanlong@amss.ac.cn (Y. Wang), zhangjun12@sjtu.edu.cn (J. Zhang), daoyidong@gmail.com (D. Dong), s.yokoyama@adfa.edu.au (S. Yokoyama), i.r.petersen@gmail.com (I.R. Petersen), h.yonezawa@adfa.edu.au (H. Yonezawa). 2022). When the operators describing a detector are diagonal in the Fock state basis, they are called phase-insensitive (otherwise phase-sensitive) detectors and can be straightforwardly identified using function fitting (Renema et al., 2012) or convex optimization (Feito et al., 2009; Lundeen et al., 2009; Natarajan et al., 2013). For phase-sensitive detectors, generally they cannot be simultaneously diagonalized and their reconstruction is thus more complicated. Existing methods include Maximum Likelihood Estimation (D'Ariano, Maccone, & Presti, 2004; Fiurášek, 2001),

Guță & Yamamoto, 2016; Nurdin & Guță, 2022). In quantum estimation and quantum system identification, a common and

essential step is to perform measurement on the quantum system

of interest. Quantum detector tomography (QDT), as the standard

technique to characterize an unknown measurement process, is

fundamental for device benchmarking and subsequent tasks such

as quantum state tomography (QST) (Hou et al., 2016; Mu, Qi,

Petersen, & Shi, 2020; Qi et al., 2013), quantum Hamiltonian

identification (Sone & Cappellaro, 2017a, 2017b; Wang et al.,

2018, 2020; Zhang & Sarovar, 2014, 2015), quantum process tomography (Fiurášek & Hradil, 2001; Wang et al., 2019; Xiao, Xue,

Dong, & Zhang, 2021) and quantum control (Dong & Petersen,



T IFA

automatica

 $[\]stackrel{\mbox{\tiny $\widehat{\mbox{$\widehat{}$}$}}}{\mbox{$\widehat{$\widehat{$}$}$}}$ This research was supported by the National Natural Science Foundation of China (62173229, 12288201), the Australian Research Council's Future Fellowship funding scheme under Project FT220100656, the Centres of Excellence under Grant CE170100012, and U.S. Office of Naval Research Global under Grant N62909-19-1-2129. The material in this paper was not presented at any conference. This paper was recommended for publication in revised form by Associate Editor Gianluigi Pillonetto under the direction of Editor Alessandro Chiuso.

^{*} Corresponding authors.

linear regression (Grandi, Zavatta, Bellini, & Paris, 2017), convexquadratic optimization (Zhang, Coldenstrodt-Ronge, et al., 2012) and Zhang, Datta, et al. (2012), and analytical two-stage solution (Wang et al., 2021). Specially, binary detectors can always be simultaneously diagonalized and thus their estimation has an analytical scheme based on Frobenius-norm projection (Wang, Dong, & Yonezawa, 2019).

For *d*-dimensional QDT, we prepare *M* different types of quantum states and the total number of copies of these states N is called resource number. Many identification algorithms assume the experimental resource is diverse enough in QDT, i.e., any detector can be uniquely determined by the measurement outcome statistics. This scenario is called informationally complete (I.C.) (Busch, 1991; Prugovečki, 1977) and the opposite scenario is called information incomplete (I.I.). In practice, the I.C. condition may not be satisfied for QDT, which results in an I.I. scenario (e.g. when $M < d^2$ for a *d*-dimensional detector). In the I.I. scenario and in certain I.C. scenarios where the probe states lie close to the I.I. set although they are still in I.C. set, the QDT problem is ill-conditioned. To solve this problem, convex optimization methods with regularization were proposed in Feito et al. (2009) and Lundeen et al. (2009) for phase-insensitive detectors and in Zhang, Coldenstrodt-Ronge, et al. (2012) and Zhang, Datta, et al. (2012) for phase-sensitive detectors. In experiments, a regularized least-square method was used in Brida, Ciavarella, Degiovanni, Genovese, Lolli, et al. (2012), Brida, Ciavarella, Degiovanni, Genovese, Migdall, et al. (2012) for phase-insensitive detectors. However, there is still a lack of closed form solutions for ODT with regularization in these existing methods. To solve this problem, we develop ODT with regularization inspired by classical transfer function identification. In the previous literature, a kernel-based regularization was proposed in Chen, Ohlsson, and Ljung (2012), Mazzoleni, Chiuso, Scandella, Formentin, and Previdi (2022), Pillonetto, Chiuso, and De Nicolao (2011), Pillonetto and De Nicolao (2010) and Pillonetto, Dinuzzo, Chen, Nicolao, and Ljung (2014), which can cope with bias-variance trade-off. For kernel-based regularization, an important problem is how to design a suitable kernel matrix. Chen (2018) and Chen et al. (2012) proposed different kernels and Chen and Andersen (2021), Chen and Ljung (2013), Mu, Chen, and Ljung (2018) and Mu, Chen, and Ljung (2021) discussed how to tune hyper-parameters in the kernel matrix and the asymptotic properties of these parameters. Further work about kernel-based regularization was studied in Chen, Andersen, Ljung, Chiuso, and Pillonetto (2014), Mu and Chen (2018), Pillonetto, Chen, Chiuso, De Nicolao, and Ljung (2022) and Pillonetto, Chen, Chiuso, Nicolao, and Ljung (2016).

In this paper, we develop regularization methods in QDT which are applicable to both phase-insensitive and phasesensitive detectors. We give a closed form solution, applicable to both the cases of I.C. and I.I. We then discuss different regularization forms and explain the advantages of using regularization in QDT. We consider no regularization as a special case. In the I.C. scenario, a common step (see e.g. Wang et al., 2019; Wang et al., 2021) is to uniformly distribute the resource for each quantum state as N/M, which is often not the optimal distribution. Without regularization, we discuss how to optimize the resource distribution for different types of probe states based on minimizing the mean squared error (MSE) of QDT. We convert this optimization problem into a semidefinite programming (SDP) problem, which can be solved efficiently. In comparison, if the resource distribution is given, the probe state design problem was discussed in Xiao, Wang, Dong, and Zhang (2022). In both the I.C. and I.I. scenarios, we also prove that under the static assumption (specific definitions in Section 4.1), the MSE scales as O(1/N) or tends to a constant, and we characterize when the MSE can reach the optimal scaling O(1/N). We propose an exact characterization of the best regularization for identifiable parameters to achieve the minimum MSE, allowing the probe states to be I.C. or I.I. In the I.C. scenario, we obtain the same best regularization form as proposed in Chen et al. (2012). We also prove the best regularization can reach the optimal scaling O(1/N) even in the I.I. scenario. Numerical examples demonstrate that the optimization of resource distribution and regularization can reduce the MSE. Then we give the reason why adaptive rank-1 regularization motivated from the best regularization fails to show an O(1/N) scaling in QDT, and we find an indication that full-rank regularization might be better. Finally, we apply our algorithm to quantum optical experiments using two-mode coherent states for binary detectors. The experimental results show that the adaptive regularization has a lower MSE compared with the Tikhonov regularization method in Wang et al. (2021). The main contributions of this paper are summarized as follows.

- (i) A closed form of regularized QDT solution is established with different regularization forms in the I.C. and I.I. scenarios. The motivations and advantages to apply regularization in QDT are discussed.
- (ii) Without regularization, we optimize the resource (probe state) distribution by converting it to a semidefinite programming (SDP) problem in the I.C. scenario.
- (iii) Under the static assumption, we prove that the MSE scales as O(1/N) or tends to a constant and we characterize when the MSE can reach the optimal scaling O(1/N). In addition, an exact characterization of the best regularization for identifiable parameters to achieve the minimum MSE is given in the I.C. and I.I. scenarios.
- (iv) Simulation results are provided to verify the effectiveness of resource distribution optimization and regularized QDT. Quantum optics experimental results are presented to demonstrate the necessity of choosing a proper regularization form to further reduce the QDT error.

This paper is organized as follows. In Section 2, we introduce the background knowledge and weighted least squares for QDT. In Section 3, we discuss different regularization forms for QDT. In Section 4, we characterize the scaling of MSE under static assumptions and the best regularization for identifiable parameters. In Section 5, we give numerical examples and in Section 6, we present experimental results. Conclusions are presented in Section 7.

Notation: For a matrix $A, A \ge 0$ means A is positive semidefinite. The conjugation and transpose (T) of A is A^{\dagger} . The trace of A is Tr(A). The rank of A is Rank(A). The identity matrix is I. The real and complex domains are \mathbb{R} and \mathbb{C} , respectively. The tensor product is \otimes . The set of all d-dimensional complex/real vectors is $\mathbb{C}^d/\mathbb{R}^d$. Row and column vectors also denoted as $\langle \psi |$ and $|\psi \rangle$, respectively. The Frobenius norm for a matrix and 2-norm for a vector are $\|\cdot\|$. The Kronecker delta function is δ . i = $\sqrt{-1}$. The diagonal matrix X formed from vector b is denoted as X = diag(b). For any $X_{d \times d} \ge 0$ with spectral decomposition $X = UPU^{\dagger}$, define \sqrt{X} or $X^{\frac{1}{2}}$ as U diag $(\sqrt{P_{11}}, \sqrt{P_{22}}, \ldots, \sqrt{P_{dd}}) U^{\dagger}$. The Pauli matrices are $\sigma_x, \sigma_y, \sigma_z$.

2. Preliminaries and identification algorithm

Here we present the background knowledge and briefly introduce the QDT identification algorithm in Wang et al. (2021). Based on this QDT identification algorithm, we introduce weighted least squares (WLS) in QDT and explain its advantages.

2.1. Quantum state and measurement

For a *d*-dimensional quantum system, its state can be described by a $d \times d$ Hermitian matrix ρ , which satisfies $\rho \ge 0$ and $\operatorname{Tr}(\rho) = 1$. When $\rho = |\psi\rangle\langle\psi|$ for some $|\psi\rangle \in \mathbb{C}^d$, we call ρ a pure state. Otherwise, ρ is called a mixed state, and can be represented using pure states $\{|\psi_i\rangle\}$ as $\rho = \sum_i c_i |\psi_i\rangle\langle\psi_i|$ where $c_i \in \mathbb{R}$ and $\sum_i c_i = 1$ with $c_i \ge 0$. A set of operators $\{P_i\}_{i=1}^n$ named Positive-operator-valued mea-

A set of operators $\{P_i\}_{i=1}^n$ named Positive-operator-valued measure (POVM) characterizes a corresponding detector as a measurement device. Each POVM element P_i is Hermitian and positive semidefinite, and together they satisfy the completeness constraint $\sum_{i=1}^n P_i = I$. When the measurements corresponding to $\{P_i\}$ are performed on ρ , the probability of obtaining the *i*th result is given by the Born's rule

$$p_i = \operatorname{Tr}\left(P_i\rho\right). \tag{1}$$

From the completeness constraint, we have $\sum_i p_i = 1$.

2.2. Problem formulation of QDT

Suppose the true values of the POVM elements are $\{P_i\}_{i=1}^n$. We design M different types of quantum states ρ_j (called probe states) and record the measurement results \hat{p}_{ij} as the estimate of $p_{ij} = \text{Tr}(P_i\rho_j)$. Each probe state has resource number N_j (i.e., N_j copies) and the total resource number is $N = \sum_{j=1}^M N_j$. Given experimental data $\{\hat{p}_{ij}\}$, the problem of QDT (Wang et al., 2021) can be formulated as

$$\min_{\{\hat{P}_i\}_{i=1}^n} \sum_{i=1}^n \sum_{j=1}^M \left[\hat{p}_{ij} - \operatorname{Tr}\left(\hat{P}_i \rho_j\right) \right]^2$$
(2)

such that $\hat{P}_i = \hat{P}_i^{\dagger}, \hat{P}_i \ge 0$ for $1 \le i \le n$ and $\sum_{i=1}^n \hat{P}_i = I$.

Let $\{\Omega_i\}_{i=1}^{d^2}$ be a complete basis set of orthonormal operators with dimension d, $\operatorname{Tr}\left(\Omega_i^{\dagger}\Omega_j\right) = \delta_{ij}$, $\Omega_i = \Omega_i^{\dagger}$ and $\operatorname{Tr}\left(\Omega_i\right) = 0$ except $\Omega_1 = I/\sqrt{d}$. Then we can parameterize the detector and probe states as $P_i = \sum_{a=1}^{d^2} \lambda_i^a \Omega_a$, $\rho_j = \sum_{b=1}^{d^2} \phi_j^b \Omega_b$. Using Born's rule, we can obtain the least squares (LS) equation

$$\hat{p}_{ij} = \phi_i^I \lambda_i + e_{ij},\tag{3}$$

where ϕ_j and λ_i are the parameterization vectors of ρ_j and P_i , respectively. Suppose the outcome for P_i of ρ_j appears n_{ij} times, and then $\hat{p}_{ij} = n_{ij}/N_j$. The estimation error is denoted as $e_{ij} = \hat{p}_{ij} - p_{ij}$. According to the central limit theorem, e_{ij} converges in distribution to a normal distribution with mean zero and variance $(p_{ij} - p_{ij}^2)/N_j$.

To propose least squares (LS) and weighted least squares (WLS) solutions in QDT, in the following we write down and solve the linear equation for each POVM element *individually*. This can be achieved by rearranging the data after implementing all the measurements. Collect the parameterization of the probe states as $X = (\phi_1, \phi_2, \dots, \phi_M)^T$. For the *i*th POVM element P_i , let $\hat{y}_i \triangleq (\hat{p}_{i1}, \hat{p}_{i2}, \dots, \hat{p}_{iM})^T$, $y_0 \triangleq ((1, \dots, 1)_{1 \times M})^T = \sum_i \hat{y}_i$, $d_{d^2 \times 1} \triangleq (\sqrt{d}, 0, \dots, 0)^T$, $e_i \triangleq (e_{i1}, \dots, e_{iM})^T$. Define $\bar{y}_i \triangleq \hat{y}_i - \frac{1}{n}y_0$ and $\theta_i \triangleq \lambda_i - \frac{1}{n}d$. Then we have

$$\bar{y_i} = X\theta_i + e_i. \tag{4}$$

Now the QDT problem can be transformed into the following form:

Problem 1. For $1 \le i \le n$, given experimental data \bar{y}_i , solve $\min_{\hat{P}_i} \|\bar{y}_i - X\theta_i\|^2$ with $\hat{P}_i \ge 0$, where $\lambda_i = \theta_i + \frac{1}{n}d$ is the parameterization of \hat{P}_i .

2.3. Weighted least squares in QDT

To solve Problem 1, the standard LS solution is

$$\hat{\theta}_{i,\mathrm{LS}} = \left(X^T X\right)^{-1} X^T \bar{y}_i,\tag{5}$$

and then the estimator for each detector is $\hat{\lambda}_{i,LS} = \hat{\theta}_{i,LS} + \frac{1}{n}d$, which is equivalent to equation (9) in Wang et al. (2021).

Although all the estimation errors e_{ij} have zero mean, they have different variances, which is called heteroscedasticity in statistics. The constrained least squares as equation (6) in Wang et al. (2021) and standard LS (5) do not consider heteroscedasticity. However, WLS consider the heteroscedasticity property and has optimal MSE. We thus consider WLS estimate

$$\hat{\theta}_{i,\text{WLS}} = \left(X^T W_i X\right)^{-1} X^T W_i \bar{y}_i,\tag{6}$$

where $W_i = \text{diag}\left(\left[\frac{N_1}{p_{i1}-p_{i1}^2}, \ldots, \frac{N_M}{p_{iM}-p_{iM}^2}\right]\right)$ is the weighting matrix. We assume that p_{ij} is not equal to 0 or 1, which is reasonable because $p_{ij} \in [0, 1]$ and generally the probability for $p_{ij} = 0$ or 1 is 0 in theory. The following are the two main advantages of using WLS:

• We can normalize the estimation errors to normal Gaussian errors and solve the heteroscedasticity problem. With increasing measurements, each e_{ij} will converge asymptotically to a Gaussian random variable with mean zero and variance $\sigma_{ij} = (p_{ij} - p_{ij}^2) / N_j$. Thus, we have $\mathbb{E}(e_i e_i^T) = W_i^{-1}$. Define $Q_i \triangleq \sqrt{W_i^{-1}} / \sigma$ for certain $\sigma > 0$. Then we multiply by Q_i^{-1} in (4) as

$$Q_i^{-1}\bar{y}_i = Q_i^{-1}X\theta_i + Q_i^{-1}e_i.$$
(7)

Let $\mathbb{E}(\cdot)$ denote the expectation with respect to all possible measurement results. The new errors have an independent identical normal distribution (i.i.d.) with

$$\mathbb{E}\left(Q_i^{-1}e_ie_i^TQ_i^{-1}\right) = \sigma^2 I.$$
(8)

Thus, all the variances of the estimation errors are normalized to σ^2 .

• For any unbiased linear estimator $\hat{\theta}_i$ for θ_i , we have (Mu et al., 2020)

$$MSEM\left(\hat{\theta}_{i,WLS}\right) = \mathbb{E}\left[\left(\hat{\theta}_{i,WLS} - \theta_{i}\right)\left(\hat{\theta}_{i,WLS} - \theta_{i}\right)^{T}\right]$$
$$= \left(X^{T}W_{i}X\right)^{-1} \leq MSEM\left(\hat{\theta}_{i}\right),$$
(9)

where $\mathsf{MSEM}\left(\cdot\right)$ is the MSE matrix. The MSE of all the POVM elements is

$$\mathbb{E}\left(\sum_{i=1}^{n} \left\|\hat{P}_{i} - P_{i}\right\|^{2}\right) = \sum_{i=1}^{n} \mathbb{E}\left(\left\|\hat{\theta}_{i} - \theta_{i}\right\|^{2}\right)$$
$$= \sum_{i=1}^{n} \operatorname{Tr}\left(\operatorname{MSEM}\left(\hat{\theta}_{i}\right)\right).$$
(10)

Hence, the WLS solution to Problem 1 has the minimum MSE.

In practice, the weighting matrix W_i is unknown and a feasible solution is to use the asymptotic estimate

$$\hat{W}_{i} = \text{diag}\left(\left[\frac{N_{1}}{\hat{p}_{i1} - \hat{p}_{i1}^{2}}, \dots, \frac{N_{M}}{\hat{p}_{iM} - \hat{p}_{iM}^{2}}\right]\right).$$
(11)

Denote $\hat{Q}_i^{-1} \triangleq \sqrt{\hat{W}_i} / \sigma$, $\tilde{y}_i \triangleq \hat{Q}_i^{-1} \bar{y}_i$, $\tilde{X}_i \triangleq \hat{Q}_i^{-1} X$, $\tilde{e}_i \triangleq \hat{Q}_i^{-1} e_i$ and the model equivalent to (4) is

$$\tilde{y}_i = \tilde{X}_i \theta_i + \tilde{e}_i, \tag{12}$$

where the variance of \tilde{e}_i is $\sigma^2 I$ and the practical asymptotic WLS (AWLS) estimate is

$$\hat{\theta}_{i,\text{AWLS}} = \left(X^T \hat{W}_i X\right)^{-1} X^T \hat{W}_i \bar{y}_i = \left(\tilde{X}_i^T \tilde{X}_i\right)^{-1} \tilde{X}_i^T \tilde{y}_i.$$
(13)

The difference between $\hat{\theta}_{i,\text{MVLS}}$ and $\hat{\theta}_{i,\text{WLS}}$ is asymptotically small in comparison with $\hat{\theta}_{i,\text{WLS}}$ (Mu et al., 2020). Thus, the estimate (13) is accurate enough and asymptotically coincides with (6). Using the LS estimate (5) or WLS estimate (13), we can obtain a POVM estimate $\hat{E}_i = \sum_{a=1}^{d^2} \left(\hat{\theta}_{i,\text{LS/WLS}} + \frac{1}{n} d \right)_a \Omega_a$ and $\mathbb{E} \| \hat{E}_i - P_i \|^2 =$ $\text{Tr} \left(\text{MSEM} \left(\hat{\theta}_i \right) \right)$. We call the error $\mathbb{E} \| \hat{E}_i - P_i \|^2$ the LS MSE for the *i*th POVM element. Note that $\{ \hat{E}_i \}_{i=1}^n$ may have negative eigenvalues due to the noise or error in the measurement results. Thus, we need further correction to obtain a positive semidefinite estimate $\{ \hat{P}_i \}_{i=1}^n$ and in this paper, we utilize the algorithm in Wang et al. (2021) to achieve this. We refer to the error $\mathbb{E} \| \hat{P}_i - P_i \|^2$ as the final MSE for the *i*th POVM element.

Remark 1. One may notice that (4) has the same linear regression form $v = X\theta + e$ as transfer function identification in system identification (Chen et al., 2012). However, there are some differences between ODT and transfer function identification for classical (non-quantum) systems. First, in QDT, more measurement data will only enhance the data accuracy in y and the dimension of y is fixed with given probe states. In transfer function identification, the dimension of y increases for more data. Second, the parameterization matrix X is determined by the given probe states and $X^T X$ can be singular (e.g., $M < d^2$) in ODT. In transfer function identification, X depends on the input data and measurement data. In practice, $X^T X$ is therefore always invertible but the condition number may be large. Thus, the standard LS cannot give an accurate estimate. Finally, the variance of the noise e is often assumed to be a constant in transfer function identification. However, in QDT, the variances of noise are usually different and decrease as O(1/N) where N is the resource number.

3. Regularization in QDT

In QDT, when the different types of probe states are similar or I.I., leading to an ill-conditioned problem, convex optimization methods with regularization were proposed in Feito et al. (2009) and Lundeen et al. (2009) for phase-insensitive detectors and in Zhang, Coldenstrodt-Ronge, et al. (2012) and Zhang, Datta, et al. (2012) for phase-sensitive ones. The motivation of introducing regularization is to mitigate the ill-conditioned property. For phase-insensitive detectors, the regularization form is chosen such that the diagonal elements of the reconstructed detector have smooth variations (Zhang, Datta, et al., 2012). However, for phase-sensitive detectors, a suitable regularization form is not easy to find. In addition, convex optimization methods cannot give a closed form solution. Therefore, in this section, we use regularization in the WLS of QDT which can give a closed form solution.

3.1. Regularized weighted least squares

In the ill-conditioned scenario, the condition number of $\tilde{X}_i^T \tilde{X}_i$ can be large or even infinite. To solve this problem, we add regularization in the weighted model (12) as

$$\left\|\tilde{y}_{i}-\tilde{X}_{i}\theta_{i}\right\|^{2}+\theta_{i}^{T}D_{i}\theta_{i},$$
(14)

where D_i is positive semi-definite and called a regularization matrix. Denote $R_i \triangleq \tilde{X}_i^T \tilde{X}_i$. After we add regularization, the estimate is changed to be

$$\hat{\theta}_{i,\text{RWLS}} = (R_i + D_i)^{-1} \tilde{X}_i^T \tilde{y}_i.$$
(15)

The expectation of $\hat{\theta}_{i,RWLS}$ is

$$\mathbb{E}\left(\hat{\theta}_{i,\text{RWLS}}\right) = \left(R_i + D_i\right)^{-1} R_i \theta_i.$$
(16)

The bias is

$$\theta_{i,\text{RWLS}}^{\text{bias}} \triangleq \mathbb{E}\left(\hat{\theta}_{i,\text{RWLS}}\right) - \theta_i = -\left(R_i + D_i\right)^{-1} D_i \theta_i.$$
(17)

Define

$$\begin{split} \tilde{\theta}_{i} &\triangleq \hat{\theta}_{i,\text{RWLS}} - \mathbb{E}\left(\hat{\theta}_{i,\text{RWLS}}\right) \\ &= (R_{i} + D_{i})^{-1} \tilde{X}_{i}^{T} \left(\tilde{y}_{i} - \tilde{X}_{i} \theta_{i}\right) \\ &= (R_{i} + D_{i})^{-1} \tilde{X}_{i}^{T} \tilde{e}_{i}, \end{split}$$
(18)

and then the MSE matrix of $\hat{\theta}_{i,RWLS}$ is

$$MSEM\left(\hat{\theta}_{i,RWLS}\right) = \mathbb{E}\left[\left(\hat{\theta}_{i,RWLS} - \theta_{i}\right)\left(\hat{\theta}_{i,RWLS} - \theta_{i}\right)^{T}\right]$$

$$= \mathbb{E}\left(\tilde{\theta}_{i}\tilde{\theta}_{i}^{T}\right) + \theta_{i,RWLS}^{\text{bias}}\left(\theta_{i,RWLS}^{\text{bias}}\right)^{T}$$

$$= (R_{i} + D_{i})^{-1}\left(\sigma^{2}R_{i} + D_{i}\theta_{i}\theta_{i}^{T}D_{i}^{T}\right)(R_{i} + D_{i})^{-1}.$$

(19)

An MSE matrix similar to (19) can be found in Chen et al. (2012) for transfer function identification with standard LS estimation. The LS MSE of QDT is Tr(MSEM) and depends on the true parameter θ_i . When the probe states are I.C., we can obtain an estimate without regularization (i.e., $D_i = 0$) and the MSE matrix becomes

$$\mathsf{MSEM}\left(\hat{\theta}_{i,\mathsf{RWLS}}\right) = \sigma^2 R_i^{-1},\tag{20}$$

which is independent of the true parameter θ_i .

Based on the development in classical system identification, several motivations of applying regularization in QDT are as follows:

- (i) Regularization is a typical solution to ill-conditioned problems. In the field of classical transfer function identification (see e.g., Chen & Ljung, 2013), the input signal is bandlimited, and then the matrix R_i may become ill-conditioned as the amount of data increases. Similarly in QDT, the input probe states can be "band-limited", in the sense that the types of the probe states are not rich enough (especially when coherent states are employed) which leads to the conversion from I.C. to I.I. This current incapability of realizing perfect number states endows R_i with a large condition number, which can be reduced by regularization while still maintaining a closed form solution.
- (ii) From an alternative point of view, regularization leverages the bias-variance trade-off. The regularization estimation is biased as (17), which can lead to an MSE smaller than that of the standard LS estimation both in the I.C. and I.I. scenarios.

There are also differences of applying regularization between QDT and classical system identification. All physical POVM elements must be positive semidefinite and sum to identify, which may affect or even guide the design of the specific regularization form in QDT. For example, Wang et al. (2021) noted that POVM elements satisfying these physical constraints always have eigenvalues in [0, 1]. Direct LS estimation for ill-conditioned QDT usually gives a large $\|\hat{\theta}_i\|$, which may have eigenvalues outside [0, 1] and become nonphysical. Therefore, the regularization $\theta_i^T D_i \theta_i$ is added to the

cost function in Wang et al. (2021) as a penalty term, promoting the satisfaction of the physical constraints. Apart from this, other differences will be detailed in Section 3.2.3.

Regularized weighted regression is also applied in quantum state tomography. For example, in Mu et al. (2020), their motivation is that the quantum state ρ is usually of low rank (Cai, Kim, Wang, Yuan, & Zhou, 2016) and thus it is reasonable to add a Tikhonov regularization as Section 3.2.2. However, in QDT, the POVM elements are not always of low rank. For example, in the continuous-variable optical experiment in the paper and in Feito et al. (2009), Lundeen et al. (2009) and Zhang, Datta, et al. (2012), the POVM elements are all full-rank. Thus we introduce and discuss more regularization forms in Section 3.2.

3.2. Different regularization forms in QDT

Here we discuss different regularization forms in QDT. Firstly, we consider no regularization (i.e., $D_i = 0$) as a special regularization form in the I.C. scenario. Since the MSE in (20) does not depend on true parameter θ_i , we propose resource distribution optimization of N_j to minimize the LS MSE with given probe states. Then we present some common regularization forms. With regularization, the LS MSE in (19) depends on true parameter θ_i . Thus we cannot optimize resource distribution as without regularization and we use a uniformly distributed $N_i = N/M$.

3.2.1. Without regularization

Without regularization, Wang et al. (2019) and Wang et al. (2021) choose $N_j = N/M$ for given probe states, which is often not the optimal distribution. Similar input design problems in classical systems and control have been widely studied and there are many existing results, e.g., D,A,E-optimal input design (Boyd & Vandenberghe, 2004). Here, we formulate and solve the problem within the framework of A-optimal design problem, where the trace of the covariance matrix (i.e., MSE) is minimized.

Let $\eta_j = \frac{N_j}{N}$, and the optimization of resource distribution problem can be formulated as

$$\min_{\{\eta_j\}_{j=1}^M} \sum_{i=1}^n \operatorname{Tr}\left(\sum_{j=1}^M \left(\eta_j w_{ij} \phi_j \phi_j^T\right)\right)^{-1}$$
s.t. $\eta_j \ge 0, \sum_{j=1}^M \eta_j = 1,$
(21)

where ϕ_j is the given parameterization vectors of ρ_j and w_{ij} is the weighted constant which we may obtain from a prior information. If we do not have a prior information, we can set $w_{ij} = 1$. This optimization problem is convex and it can be converted to a semidefinite programming (SDP) problem

$$\min_{\{\eta_j\}_{j=1}^{M}, \{u_k\}_{k=1}^{d^2}} \sum_{k=1}^{d^2} u_k \\
\text{s.t.} \left[\sum_{j=1}^{M} \eta_j w_{ij} \phi_j \phi_j^T \quad v_k \\
v_k^T \quad u_k \right] \ge 0, \\
1 \le k \le d^2, \ 1 \le i \le n, \\
\eta_j \ge 0, \ \sum_{j=1}^{M} \eta_j = 1,$$
(22)

where v_k is the *k*th unit vector. Using CVX (Grant & Boyd, 2008, 2014), we can solve (22) efficiently. Note that $N_j = \eta_j N$ may not be an integer, and we need to round it up or down. In comparison, if the resource distribution is given, the probe state design problem was discussed in Xiao et al. (2022) based on minimizing an upper bound on the MSE and the condition number.

3.2.2. Tikhonov regularization

A most common regularization form is in a Tikhonov sense (Boyd & Vandenberghe, 2004). In QDT, a natural method is to choose regularization matrix as

$$D_i^{\text{Tikhonov}} = cI, \tag{23}$$

where *c* is a positive constant. Wang et al. (2021) did not use WLS and chose $D_i = \frac{c}{N}I$ which is Tikhonov regularization, because

$$\hat{\theta}_{i,\text{RWLS}} = \left(X^T X + \frac{c}{N}I\right)^{-1} X^T \bar{y}_i$$

$$= \left(X^T N I X + cI\right)^{-1} X^T N I \bar{y}_i,$$
(24)

where the weighted matrix is NI instead of (11).

3.2.3. Kernel-based regularization

In transfer function identification, Chen et al. (2012), Pillonetto et al. (2011), Pillonetto and De Nicolao (2010) and Pillonetto et al. (2014) proposed kernel-based regularization and explained regularization in a Bayesian perspective. We assume the true parameter θ_i is a random variable and has a Gaussian distribution with zero mean and covariance matrix S_i :

$$\theta_i \sim \mathcal{N}(\mathbf{0}, \mathbf{S}_i)$$
 (25)

Therefore, the posterior estimate is

$$\hat{\theta}_{i}^{\text{post}} = \left(S_{i}R_{i} + \sigma^{2}I\right)^{-1}S_{i}F_{i} = \left(R_{i} + \sigma^{2}S_{i}^{-1}\right)^{-1}F_{i},$$
(26)

where $F_i \triangleq \tilde{X}_i^T \tilde{y}_i$. If S_i is singular, we can use the first equality of (26) to obtain the estimate. This posterior estimate is the same as the regularized estimate if the regularization matrix D_i is chosen as (Chen et al., 2012)

$$D_i = \sigma^2 S_i^{-1}. \tag{27}$$

This gives an insight into how to choose the regularization matrix D_i or kernel matrix S_i : Let it reflect the correlations of the parameters (Chen et al., 2012).

To use the kernel-based regularization in QDT, we need to solve two problems

- (i) In the Bayesian perspective for kernel-based regularization, the mean of the unknown parameters is zero. But in QDT, the mean of the unknown parameters λ_i is usually not zero.
- (ii) Heteroscedasticity: In transfer function identification, it is usually assumed that the noises have the same variances. But the estimation errors e_{ij} usually have different variances in QDT.

The first problem is solved by modeling in (4) where the unknown parameter θ_i becomes zero-mean. For the second problem, WLS (12) solves the heteroscedasticity problem.

There are two advantages of using kernel-based regularization in QDT compared with using kernel-based regularization in transfer function identification:

- (i) In transfer function identification, we need to identify the variance of the noise firstly, while we already know the approximate variance of the estimation error in QDT from measurement data.
- (ii) In transfer function identification, the problem dimension increases as more data are generated, resulting in increased difficulty. While in QDT, more data will only enhance the data accuracy and the dimension is fixed with given probe states.

One limit using kernel-based regularization in QDT is that without prior knowledge the parameter θ_i does not have the property of impulse responses of transfer functions which usually

decay exponentially (Chen et al., 2012). In this paper, we mainly choose DI kernel which only represents the auto-correlation for each coefficient of QDT as

$$S_i^{\text{DI}}(k,j) = \begin{cases} c\mu^k, & \text{if } k = j, \\ 0, & \text{otherwise,} \end{cases}$$
(28)

where $c \ge 0$, $0 \le \mu \le 1$. If we have more prior knowledge such as the correlation between different coefficients, we can design more suitable kernels as in transfer function identification. For example, when the detector is close to a phase-insensitive detector, i.e., the POVM elements are close to diagonal matrices in the Fock state basis, the true value θ_i is close to sparse, which is similar to the decay behavior of impulse responses for stable transfer functions in system identification. Therefore, we can apply TC and DC kernels (Chen & Ljung, 2013; Chen et al., 2012) in transfer function identification

$$S_i^{\text{TC}}(k,j) = c \min\left(\mu^j, \mu^k\right),\tag{29}$$

where $c \ge 0$, $0 \le \mu \le 1$ and

$$S_i^{\rm DC}(k,j) = c \mu_1^{|k-j|} \mu_2^{(k+j)/2},$$
(30)

where $c \ge 0$, $-1 \le \mu_1 \le 1$ and $0 \le \mu_2 \le 1$.

3.2.4. Best regularization (in the I. C. scenario)

For true parameter θ_i , two natural questions are whether there exists an optimal regularization matrix and if there exists an optimal regularization matrix, does it depend on θ_i ? Chen et al. (2012) has discussed these problems in transfer function identification and the result also holds for QDT. The MSE matrix in (19) can be rewritten using S_i as

$$MSEM\left(\hat{\theta}_{i,RWLS}\right) = \left(S_{i}R_{i} + \sigma^{2}I\right)^{-1} \left(\sigma^{2}S_{i}R_{i}S_{i} + \sigma^{4}\theta_{i}\theta_{i}^{T}\right) \left(R_{i}S_{i} + \sigma^{2}I\right)^{-1}.$$
(31)

When R_i is invertible, the following matrix inequality (Chen et al., 2012; Eldar, 2006)

$$\left. \mathsf{MSEM}\left(\hat{\theta}_{i,\mathsf{RWLS}}\right) \right|_{S_{i}=K} \ge \left. \mathsf{MSEM}\left(\hat{\theta}_{i,\mathsf{RWLS}}\right) \right|_{S_{i}=\theta_{i}\theta_{i}^{T}}$$
(32)

holds for any $K \ge 0$. Later, in Theorem 2, we will extend this inequality to the case where R_i is singular. Thus, ideally the best choice of regularization always includes

$$S_i^{\text{best}} = \theta_i \theta_i^T, \tag{33}$$

which yields the corresponding optimal regularized estimate

$$\hat{\theta}_i^{\text{best}} = \left(\theta_i \theta_i^T R_i + \sigma^2 I\right)^{-1} \theta_i \theta_i^T F_i, \tag{34}$$

with $R_i = \tilde{X}_i^T \tilde{X}_i$ and $F_i = \tilde{X}_i^T \tilde{y}_i$. The theoretically best regularization depends on the unknown parameter and cannot be used in practice.

A natural question is that, is $\theta_i \theta_i^T$ the only choice for S_i to result in the best regularization? Eldar (2006) has given a positive answer for the I.C. scenario. For the I.I. scenario we will give a negative answer in Section 4.2.

3.2.5. Adaptive regularization

As motivated by the best regularization, we can propose adaptive regularization with rank-1 kernel matrix which is similar to the rank-1 kernel matrix for transfer function identification in Chen et al. (2014). Firstly, we consider a two-step adaptive regularization. In the first step, we use Tikhonov or kernel-based regularization and we can obtain a rough estimate $\hat{\theta}_i^0$ with certain kernel matrix $S_i^{(1)}$. Then in the second step, we repeat using the measurement data in the first step, but now the regularization matrix is adaptively chosen as

$$S_i^{\text{rank-1}} = \hat{\theta}_i^0 \left(\hat{\theta}_i^0 \right)^T.$$
(35)

The following analysis and Theorem 1 in the next section indicate that full-rank kernel matrix may be better than rank-1 kernel matrix, because a full-rank S_i does not induce a dimension reduction from R(B) to $R(S_iB)$. Therefore, we also consider to use full-rank kernel matrix as

$$S_i^{\text{full-rank}} = S_i^{\text{rank-1}} + S_i^{\text{DI/TC/DC}},$$
(36)

in Section 5.

It is an important problem to determine the kernel matrix and some different kernels are proposed in transfer function identification. For a structure-given kernel matrix, optimization of the hyper-parameters (such as c, μ in (28)) in the kernel matrix has been discussed in Chen (2018), Chen et al. (2014) and Chen and Ljung (2013). However, the question of how to choose the optimal adaptive kernel matrix is still an open problem.

4. Characterizing the MSE of QDT with regularization

4.1. On the MSE scaling

To analyze the performance of different regularization methods, we characterize the asymptotic behavior of the estimation error, e.g., MSE. Without loss of generality, we can always normalize the variances of the estimation errors to 1, i.e., $\sigma^2 = 1$ in (8). We give the following assumptions.

Assumption 1. The probe state parameterization matrix *X* is given. The kernel matrix S_i is given. For each $1 \le j \le n$, $\lim_{N\to\infty} \frac{N_j}{N} = h(j)$ where h(j) is a constant in [0, 1] depending on *j*.

We refer to Assumption 1 as the *static assumption*. With Assumption 1, the probe state parameterization matrix and kernel matrix are given as constant matrices which do not change in our analysis and the resource distribution for each probe state can change as *N* increases. But the limit of the ratio is a constant and can be 0 or 1. We say that the random sequence $\{\xi_N\}$ converges almost surely to a random variable ξ if $P(\lim_{N\to\infty} ||\xi_N - \xi||_2 = 0) = 1$, which can be written as $\xi_N \xrightarrow{a.s.} \xi$ as $N \to \infty$. For the weighted matrix \hat{W}_i , its deviation from the true value W_i has been derived in Mu et al. (2020) as

$$\hat{W}_{i} = \operatorname{diag}\left(\left[\frac{N_{1}}{\hat{p}_{i1} - \hat{p}_{i1}^{2}}, \dots, \frac{N_{M}}{\hat{p}_{iM} - \hat{p}_{iM}^{2}}\right]\right)$$
$$= \left(1 + O\left(\frac{1}{\sqrt{N}}\right)\right) W_{i}.$$
(37)

We define

$$B \triangleq \lim_{N \to \infty} \frac{X^T W_i X}{N}, \hat{B}_N \triangleq \frac{X^T \hat{W}_i X}{N},$$
(38)

where the normalized weighted parameterization matrix $\hat{B}_N = \left(1 + O\left(1/\sqrt{N}\right)\right) B$ for constant matrix B because $\lim_{N\to\infty} \frac{N_j}{N} =$ constant. Therefore, $\hat{B}_N \xrightarrow{a.s.} B$ as $N \to \infty$.

We denote R(X) as the range space of X and N(X) as the null space of X. Then we propose the following theorem to characterize the MSE.

Theorem 1. In the regularization-based QDT, if the ith POVM element satisfies the static assumption, then its LS MSE $\mathbb{E} \| \hat{E}_i - P_i \|^2$

and final MSE $\mathbb{E} \|\hat{P}_i - P_i\|^2$ both scale as O(1/N) if and only if the true values of the unknown parameters satisfy $\theta_i \in R(S_iB)$. Otherwise, the LS MSE $\mathbb{E} \|\hat{E}_i - P_i\|^2$ converges to a positive value.

Proof. For the *i*th POVM element, according to (31) and $\sigma^2 = 1$, the MSE is

$$Tr \left[MSEM |_{S_i} \right]$$

= $Tr \left[(S_iR_i + I)^{-1} (S_iR_iS_i + \theta_i\theta_i^T) (R_iS_i + I)^{-1} \right]$
= $Tr \left\{ [(S_iR_i + I) (R_iS_i + I)]^{-1} (S_iR_iS_i + \theta_i\theta_i^T) \right\},$ (39)

where $R_i = X^T \hat{W}_i X$. We define

$$A_{1} \triangleq (S_{i}R_{i} + I) (R_{i}S_{i} + I)$$

= $\left(NS_{i}\hat{B}_{N} + I\right) \left(N\hat{B}_{N}S_{i} + I\right),$ (40)

and

$$A_2 \triangleq S_i R_i S_i = N S_i \hat{B}_N S_i. \tag{41}$$

Now the MSE becomes Tr $(A_1^{-1}(A_2 + \theta_i \theta_i^T))$. We then introduce the following lemma

Lemma 1 (*Cui, Li, & Sze, 2017; Wu, 1988*). For an $n \times n$ complex matrix *T*, the following statements are equivalent:

(1) T = AB, where $A, B \ge 0$;

- (2) T = AB, where A > 0 and $B \ge 0$;
- (3) T is similar to a nonnegative diagonal matrix.

From Lemma 1, S_iB is similar to a nonnegative diagonal matrix and we assume $S_iB = Q^{-1}\Sigma_1Q$ where $\Sigma_1 = \text{diag}(\Sigma_{11}, \Sigma_{12})$ and Σ_{11} is a $k \times k$ positive diagonal matrix, Σ_{12} is a $(d^2 - k) \times (d^2 - k)$ zero matrix. Therefore, $NS_iB + I$ can also be diagonalized by Q as

$$NS_{i}B + I = Q^{-1} \operatorname{diag} \left(\left[\tau_{1}, \dots, \tau_{d^{2}} \right] \right) Q$$

= Q^{-1} \operatorname{diag} \left(N \Sigma_{11} + I_{k}, I_{d^{2}-k} \right) Q, (42)

where $\tau_1 \ge \cdots \ge \tau_{d^2} > 0$, $\tau_j = O(N)$ for $1 \le j \le k$ and $\tau_j = 1$ for $k+1 \le j \le d^2$ and the corresponding eigenvectors are $\{u_j\}_{j=1}^{d^2}$. As $N \to \infty$, we have

$$\lim_{N \to \infty} (NS_i B + I)^{-1} = \lim_{N \to \infty} Q^{-1} \operatorname{diag} (N \Sigma_{11} + I_k, I_{d^2 - k})^{-1} Q$$
(43)
= $Q^{-1} \operatorname{diag} (0, I_{d^2 - k}) Q,$

and thus $(NS_iB + I)^{-1}$ tends to a constant matrix. Since

$$I - (NS_iB + I)^{-1} = (NS_iB + I)^{-1}NS_iB,$$

it is also a bounded matrix and tends to a constant matrix as $N \rightarrow \infty$. Let the spectral decomposition of *B* be

$$B = V \Sigma_2 V^T = V \operatorname{diag} (\Sigma_{21}, 0) V^T.$$
(44)

Thus, the Moore–Penrose inverse of *B* is

$$\tilde{B} = V \operatorname{diag}\left(\Sigma_{21}^{-1}, 0\right) V^{T}, \tag{45}$$

which is a constant matrix and BBB = B. Therefore, the first term of MSE is

$$\operatorname{Tr} \left(A_{1}^{-1} A_{2} \right)$$

$$\stackrel{a.s.}{\to} \operatorname{Tr} \left((NBS_{i} + I)^{-1} (NS_{i}B + I)^{-1} NS_{i}BS_{i} \right)$$

$$= \frac{1}{N} \operatorname{Tr} \left((NS_{i}B + I)^{-1} NS_{i}B \cdot \tilde{B} \cdot NBS_{i} (NBS_{i} + I)^{-1} \right)$$

$$= O\left(\frac{1}{N}\right),$$
(6)

because the term $Tr(\cdot)$ is bounded and tends to a constant. Therefore, the first term of MSE always scales as $O(\frac{1}{N})$. Then we discuss the scaling of the second part of MSE

$$\operatorname{Tr}\left(A_{1}^{-1}\theta_{i}\theta_{i}^{T}\right) \xrightarrow{a.s.} \theta_{i}^{T} \left(NBS_{i}+I\right)^{-1} \left(NS_{i}B+I\right)^{-1}\theta_{i}.$$
(47)

If θ_i is a linear combination of u_j for $1 \le j \le k$, we have

$$\operatorname{Tr}\left(A_{1}^{-1}\theta_{i}\theta_{i}^{T}\right)=O\left(\frac{1}{N^{2}}\right).$$
(48)

Otherwise, if θ_i is not a linear combination of u_j , $1 \le j \le k$, Tr $(A_1^{-1}\theta_i\theta_i^T)$ tends to a positive number independent of *N*.

Therefore, for the LS MSE $\mathbb{E}\|\hat{E}_i - P_i\|^2$, it scales as O(1/N) if and only if the true parameter θ_i is a linear combination of u_j for $1 \le j \le k$, i.e., $\theta_i \in R(S_iB)$. Since $\{\hat{E}_i\}_{i=1}^n$ may have negative eigenvalues, we use the algorithm in Wang et al. (2021) to further obtain a positive semidefinite estimate $\{\hat{P}_i\}_{i=1}^n$. The error analysis in Wang et al. (2021) has shown that

$$\sum_{i=1}^{n} \left\| \hat{P}_{i} - P_{i} \right\|^{2} = (dn + 2\sqrt{d}n + 1)O\left(\sum_{i=1}^{n} \left\| \hat{E}_{i} - P_{i} \right\|^{2} \right).$$
(49)

Therefore, if $\mathbb{E}\sum_{i=1}^{n} \left\| \hat{E}_{i} - P_{i} \right\|^{2} = O(1/N)$, we have $\mathbb{E}\sum_{i=1}^{n} \left\| \hat{P}_{i} - P_{i} \right\|^{2} = O(1/N)$ and thus the final MSE $\mathbb{E} \left\| \hat{P}_{i} - P_{i} \right\|^{2}$ also scales as O(1/N). Using (42) and (47), if the true parameter θ_{i} is not the linear combination of u_{j} for $1 \leq j \leq k$, i.e., $\theta_{i} \notin R(S_{i}B)$, the LS MSE $\mathbb{E} \left\| \hat{E}_{i} - P_{i} \right\|^{2}$ tends to a positive value. \Box

Remark 2. In Theorem 1, when $\theta_i \notin R(S_iB)$, the behavior of the final MSE $\mathbb{E} \| \hat{P}_i - P_i \|^2$ is still difficult to characterize. This problem does not exist for a full-rank detector when the resource number *N* is large enough, because the LS or WLS estimate already satisfies the positive semidefine constraint and we do not need to correct \hat{E}_i .

Note that when S_iB is full-rank, i.e., S_i and B are both positive definite, the condition $\theta_i \in R(S_iB)$ is always satisfied. Therefore, the MSE always scales as O(1/N). Thus, when the types of different probe states are I.C., for any positive definite kernel matrix S_i , the MSE always scales as O(1/N). However, when the probe states are I.I., the condition $\theta_i \in R(S_iB)$ is difficult to be satisfied in practice. Thus, without special prior knowledge, for almost all regularization forms, the LS MSE will tend to a constant when N tends to infinity. In addition, as M decreases, for given S_i , this condition may become more difficult to be satisfied because $R(S_iB)$ may become smaller. Thus, rank-1 adaptive regularization as in (35) is not a good choice and full-rank kernel matrix as in (36) may be better. The above analysis can help understand the boundary of the ability of employing regularization in QDT.

Remark 3. A similar problem was also discussed as Theorem 2.1 in Chen et al. (2014) for transfer function identification in the I.C. scenario. There a condition to realize unbiased estimation of the true parameters with regularization was given. Here, by allowing the probe states to be I.C. or I.I., we give a stronger result about the scaling of LS MSE as O(1/N) or tends to a constant for QDT. Our result can also be applied to the case when the variance of noise scales as O(1/N), which is typical in the scenario where only statistical noise is considered in quantum measurement.

4.2. On the best regularization allowing I. I.

We now consider the best regularization which has minimum MSE. It is given by (33) in the I.C. scenario. Here we aim to

46)

characterize the I.I. case. From (42) we know $NS_iB + I$ is always invertible. Define $L_i \triangleq -(NS_iB + I)^{-1}$ and thus

$$I + L_i = (NS_iB + I)^{-1} NS_iB = -NL_iS_iB.$$
(50)

Therefore, we have

$$(I+L_i)\tilde{B} = -NL_i S_i B\tilde{B}.$$
(51)

Then we propose the following theorem to characterize the best kernel matrix, allowing *B* to be singular.

Theorem 2. For the ith POVM element with true parameter θ_i and normalized weighted parameterization matrix *B* as (44), define $\Gamma \triangleq \{M \mid M = \theta_i \theta_i^T + V \operatorname{diag}(0, Z_3) V^T, Z_3 \ge 0, \operatorname{dim}(Z_3) = d^2 - \operatorname{rank}(B)\}$. If $\theta_i \in R(B)$, then S_i achieves the minimum of the LS MSE $\mathbb{E} \|\hat{E}_i - P_i\|^2$ (i.e., S_i is the best regularization) if and only if $S_i \in \Gamma$.

Proof. For the LS MSE $\|\hat{E}_i - P_i\|^2$ with kernel matrix S_i , using (40) and (41), it can be rewritten as

$$\operatorname{Tr}\left[\operatorname{MSEM}|_{S_{i}}\right] = \operatorname{Tr}\left[\left(NS_{i}B+I\right)^{-1}\left(NS_{i}BS_{i}+\theta_{i}\theta_{i}^{T}\right)\left(NBS_{i}+I\right)^{-1}\right] = \operatorname{Tr}\left[\frac{\left(I+L_{i}\right)\tilde{B}\left(I+L_{i}\right)^{T}}{N}+L_{i}\theta_{i}\theta_{i}^{T}L_{i}^{T}\right].$$
(52)

where *B* is defined in (45). Define $g(L_i)$ to be the last line of (52). Since $g(L_i)$ is convex in L_i , we can find the minimum value by setting the derivative to be zero as

$$\frac{dg}{dL_i} = \frac{2\ddot{B} + 2L_i\ddot{B}}{N} + 2L_i\theta_i\theta_i^T = 0.$$
(53)

If there exists $S_i \ge 0$ so that (53) holds for the corresponding $L_i|_{S_i}$, then such an S_i is the optimal solution to minimize the MSE (52). We tentatively plug S_i in (53), which (using (51)) becomes $2L_i\left(-S_iB\tilde{B} + \theta_i\theta_i^T\right) = 0$, equivalent to

 $\tilde{B}BS_i = \theta_i \theta_i^T.$ (54)

Since $\theta_i \in R(B)$, we let $\theta_i = Bb$ and then (54) becomes

$$V\begin{bmatrix} I_{21} & 0\\ 0 & 0 \end{bmatrix} V^{T}S_{i} = V\begin{bmatrix} \Sigma_{21} & 0\\ 0 & 0 \end{bmatrix} V^{T}bb^{T}V\begin{bmatrix} \Sigma_{21} & 0\\ 0 & 0 \end{bmatrix} V^{T}.$$
(55)

Denote $V^T b = \begin{bmatrix} p \\ q \end{bmatrix}$, $V^T S_i V = \begin{bmatrix} Z_1 & Z_2 \\ Z_2^T & Z_3 \end{bmatrix}$. Then (55) can be simplified as

$$\begin{bmatrix} I_{21} & 0\\ 0 & 0 \end{bmatrix} \begin{bmatrix} Z_1 & Z_2\\ Z_2^T & Z_3 \end{bmatrix} = \begin{bmatrix} Z_1 & Z_2\\ 0 & 0 \end{bmatrix}$$
$$= \begin{bmatrix} \Sigma_{21} & 0\\ 0 & 0 \end{bmatrix} \begin{bmatrix} p\\ q \end{bmatrix} \begin{bmatrix} p^T & q^T \end{bmatrix} \begin{bmatrix} \Sigma_{21} & 0\\ 0 & 0 \end{bmatrix}$$
$$= \begin{bmatrix} \Sigma_{21} p p^T \Sigma_{21} & 0\\ 0 & 0 \end{bmatrix},$$
(56)

and thus

$$Z_1 = \Sigma_{21} p p^T \Sigma_{21}, Z_2 = 0.$$
Since
(57)

$$\theta_i \theta_i^T = V \begin{bmatrix} \Sigma_{21} p p^T \Sigma_{21} & 0\\ 0 & 0 \end{bmatrix} V^T,$$
(58)

all solutions to (54) can be expressed as

$$S_i = V \begin{bmatrix} \Sigma_{21} p p^T \Sigma_{21} & 0\\ 0 & Z_3 \end{bmatrix} V^T = \theta_i \theta_i^T + V \begin{bmatrix} 0 & 0\\ 0 & Z_3 \end{bmatrix} V^T, \quad (59)$$

where Z_3 is positive semidefinite. Therefore, the solution set of (54) is exactly characterized by Γ where

$$\Gamma \triangleq \left\{ \begin{array}{l} M \mid M = \theta_i \theta_i^T + V \operatorname{diag}\left(0, Z_3\right) V^T, Z_3 \ge 0, \\ \operatorname{dim}(Z_3) = d^2 - \operatorname{rank}(B) \end{array} \right\}.$$
(60)

For all the best regularizations S_i in Γ , we have $S_iB = \theta_i\theta_i^T B$. This gives the minimum value of the MSE, which can be calculated as

$$\operatorname{Tr}\left(\operatorname{MSEM}|_{S_{i}\in\Gamma}\right) = \operatorname{Tr}\left(\operatorname{MSEM}|_{\theta_{i}\theta_{i}^{T}}\right)$$
$$= \operatorname{Tr}\left[\left(N\theta_{i}\theta_{i}^{T}B+I\right)^{-1}\left(N\theta_{i}\theta_{i}^{T}B\theta_{i}\theta_{i}^{T}\right)\right.$$
$$\left.+\left.\theta_{i}\theta_{i}^{T}\right)\left(NB\theta_{i}\theta_{i}^{T}+I\right)^{-1}\right]$$
$$= \operatorname{Tr}\left[\theta_{i}\theta_{i}^{T}\left(NB\theta_{i}\theta_{i}^{T}+I\right)^{-1}\right].$$
(61)

Remark 4. Note that the best regularization can minimize $\mathbb{E} \|\hat{E}_i - P_i\|^2$ instead of $\mathbb{E} \|\hat{P}_i - P_i\|^2$. The question to choose the best regularization to minimize the final MSE $\mathbb{E} \|\hat{P}_i - P_i\|^2$ where $\hat{P}_i \ge 0$ is still an open problem. Moreover, in practice, we do not know the true values of *B* and θ_i . One possible solution is to use a rough estimate $\hat{\theta}_i$ and \hat{B}_N to replace θ_i and *B* in Γ . In this case, there may exist an optimal choice of Z_3 to achieve the minimum MSE and we leave it as an open problem.

Here, we compare Theorems 1 and 2. If $\theta_i \in R(B)$, then for any full-rank kernel matrix S_i , $\theta_i \in R(S_iB)$ and thus the MSE scales as O(1/N). For any $S_i \in \Gamma$, we can obtain the minimum MSE. In addition, $\theta_i \in R(\theta_i \theta_i^T B) = R(S_iB)$, and thus the MSE also scales as O(1/N). If $\theta_i \in N(B)$, all the ideal measurement data p_{ij} are zero, i.e., we cannot obtain any information from the measurement data. Therefore, θ_i is not identifiable. If $\theta_i = \theta_{i,1} + \theta_{i,2}$ where $\theta_{i,1} \neq 0$, $\theta_{i,1} \in R(B)$ and $\theta_{i,2} \neq 0$, $\theta_{i,2} \in N(B)$, then $\theta_{i,1}$ is identifiable and $\theta_{i,2}$ is not identifiable. Therefore, we only aim to identify $\theta_{i,1}$ and the discussion is the same as $\theta_i \in R(B)$.

We then consider two special cases. The first one is that *B* is full-rank. Therefore, $\theta_i \in R(B)$ is always satisfied and the unique best kernel matrix is $S_i = \theta_i \theta_i^T$ which is the same as Chen et al. (2012). The second one is $S_i = \gamma \theta_i \theta_i^T$ where γ is a positive constant. Even if $\theta_i \notin R(B)$, we still have $\theta_i \in R(S_iB)$ ($\theta_i \notin N(B)$, otherwise $R(S_iB) = 0$), thus the MSE also scales as O(1/N). Note that all the above discussion is based on the assumption that *N* tends to infinity. When *N* is small, the performance of the regularization forms will be shown through simulation in Section 5.

5. Numerical simulation

In this section, the evaluation index is the sum of final MSEs $\mathbb{E} \sum_{i=1}^{n} \left\| \hat{P}_i - P_i \right\|^2$ and we discuss two commonly used classes of probe states for QDT. The first one involves *d* dimensional pure states $\rho = |\psi\rangle\langle\psi|$ where $|\psi\rangle$ is the superposition of *d* dimensional Fock states as

$$\psi\rangle = \sum_{i=1}^{d} c_i |i\rangle.$$
(62)

In Xiao et al. (2022), an analysis indicates that pure states may perform better than mixed states for QDT to minimize MSE.

Another class of commonly used probe states for QDT is the coherent states, because they are more straightforward to be prepared. A coherent state is denoted as $|\alpha\rangle$ where $\alpha \in \mathbb{C}$ and it can be expanded using Fock states as

$$|\alpha\rangle = e^{-\frac{|\alpha|^2}{2}} \sum_{i=0}^{\infty} \frac{\alpha^i}{\sqrt{i!}} |i\rangle.$$
(63)

Coherent states are in essence infinite dimensional. Denote the corresponding *d*-dimensional truncation as

$$|\alpha_d\rangle \triangleq \mathrm{e}^{-rac{|lpha|^2}{2}} \sum_{i=0}^{d-1} rac{lpha^i}{\sqrt{i!}} |i\rangle.$$

To estimate a *d* dimensional detector, in the simulation we assume that the outcomes generated by the residual signal $\text{Tr}[(|\alpha\rangle - |\alpha_d\rangle) (\langle \alpha | - \langle \alpha_d |)]$ are all included in the outcomes of the last POVM element. Since we truncate the coherent state in *d*-dimension, $\text{Tr}(|\alpha_d\rangle \langle \alpha_d |) < 1$ but for pure states in (62) $\text{Tr}(\rho) = 1$. Here we discuss resource distribution optimization without regularization and different regularization forms under the uniformly distributed resources.

When applying kernel-based regularization, an important problem is to determine the hyper-parameters (such as c, μ in (28), (29) and (30)) in the kernel matrix S_i . In this paper, we apply the same kernel matrix for all the POVM elements and use cross-validation in Chen et al. (2012) to determine these hyper-parameters:

- (i) Split the probe states randomly into two parts: an estimation data part with probe state parameterization matrix X_1 and a validation data part with probe state parameterization matrix X_2 .
- (ii) Collect all the hyper-parameters in a vector ω . Then estimate the detector as $\bar{\theta}_i$ using the measurement data from X_1 for different candidate values of hyper-parameters $\omega \in \bar{\Omega}$ where $\bar{\Omega}$ is a finite set in our paper.
- (iii) Using the validation data from X_2 , we find

$$\omega_0 = \arg\min_{\omega\in\bar{\Omega}} \sum_{i=1}^n \|\hat{y}_i - X_2\bar{\theta}_i(\omega)\|^2.$$
(64)

The model can then be re-estimated for this ω_0 using all the probe states. Other methods to determine the hyperparameters can also be found in Chen and Ljung (2013) and Chen et al. (2012).

5.1. Superposed Fock states

We consider a 4 dimensional three-valued phase-sensitive detector, which is close to phase-insensitive detector as

$$\begin{split} P_1^{(4)} &= \begin{bmatrix} 0.1 & 0 & 0.002 - 0.005i & 0.003 + 0.007i \\ 0 & 0.2 & 0 & 0 \\ 0.002 + 0.005i & 0 & 0.3 & 0 \\ 0.003 - 0.007i & 0 & 0 & 0.4 \end{bmatrix}, \\ P_2^{(4)} &= \begin{bmatrix} 0.2 & 0.001 + 0.002i & 0 & 0 \\ 0.001 - 0.002i & 0.2 & 0 & 0 \\ 0 & 0 & 0.3 & 0 \\ 0 & 0 & 0 & 0.4 \end{bmatrix}, \\ P_3^{(4)} &= I - P_1^{(4)} - P_2^{(4)}. \end{split}$$

Using the algorithm in Johnston (2016) and Miszczak (2012), we generate 20 different types of 4 dimensional pure states. To determine the hyper-parameters in the DI kernel, we use 16 pure states as estimation data and 4 pure states as validation data. We use different regularization forms including no regularization ((23) with c = 0), Tikhonov regularization ((23) with c = 10),



Fig. 1. The error scalings of different regularization forms with WLS using 20 types of 4 dimensional pure states. When the resource number $N > 10^6$, all the MSEs scale as O(1/N) satisfying Theorem 1.



Fig. 2. The MSE comparison between average and optimized resource distribution using 20 types of 4 dimensional pure states.

kernel-based regularization ((28) with c = 0.1, $\mu = 0.9$), rank-1 adaptive regularization, full-rank adaptive regularization (see Section 4) and the best regularization (33). The best regularization is the lower bound of MSE and depends on true value of θ . Therefore, it cannot be used in practice and we aim to achieve regularization closest to the best regularization. For rank-1 adaptive regularization, we use kernel-based regularization ((28) with c = 0.1, $\mu = 0.9$) in step 1 and (35) in step 2. For full-rank adaptive regularization, we use kernel-based regularization ((28) with c = 0.1, $\mu = 0.9$) in step 1 and (36) in step 2. For each resource number, we run the algorithm 100 times and obtain the average MSE and standard deviation.

The results are shown in Fig. 1. The best regularization scales as O(1/N) satisfying Theorem 1. When the resource number N < N10⁶, the MSEs of kernel-based regularization and adaptive regularization are a little smaller than Tikhonov regularization and no regularization. In addition, full-rank adaptive regularization has a little smaller MSE than rank-1 adaptive regularization. When the resource number $N > 10^6$, all the MSEs scale as O(1/N)satisfying Theorem 1. Since these 4 dimensional pure states are I.C., without regularization, we also consider resource distribution optimization. We compare the MSE of the case with averagely distributed resources N/M ("Average" in Fig. 2) and the MSE of the case with optimized resource distribution ("Optimized" in Fig. 2) by solving (22). For each resource number N, we run the algorithm 100 times and obtain the average MSE and standard deviation. The results are shown in Fig. 2. We can obtain a lower MSE with resource distribution optimization and both MSEs scale as O(1/N) when $N > 10^5$.



Fig. 3. The MSE comparison of DI, TC and DC kernels with WLS using 10 types of 4 dimensional pure states. All the MSEs tend to constants as predicted by Theorem 1 and Remark 2 because $\theta_i \in R(S_iB)$ does not hold.

Then we generate only 10 random types of 4 dimensional pure states. To determine the hyper-parameters in the different kernels, we use 8 pure states as the estimation data and 2 pure states as the validation data. Here we assume that we have the prior knowledge that the true detector is close to a phase-insensitive detector. We choose Pauli basis $\Omega = \{I \otimes I, I \otimes \sigma_z, \sigma_z \otimes I, I \otimes \sigma_z, \sigma_z \otimes I\}$ $\sigma_z \otimes \sigma_z, \ldots, \sigma_x \otimes \sigma_y \} /2$. Thus, the absolute values of the first four elements in θ_i are significantly larger than zero and all the other values in θ_i are close to zero, which is similar to the impulse responses of stable transfer functions in the system identification. Therefore, we use DI kernel ((28) with $c = 0.1, \mu = 0.9$), TC kernel ((29) with $c = 0.9, \mu = 0.8$) and DC kernel ((30) with $c = 0.1, \mu_1 = 0.2, \mu_2 = 0.9$), and compare their performance. The results are shown in Fig. 3. Compared with DI kernel, the MSE of DC kernel is improved by 28.3% when $N > 10^6$, which indicates that DC kernel is suitable to be applied for calibrating phase-insensitive detectors.

Hence we change DI kernel to DC kernel ((30) with $c = 0.1, \mu_1 = 0.2, \mu_2 = 0.9$) in this case and the results are shown in Fig. 4. In this I.I. scenario, there does not exist a unique solution for WLS (13) without regularization. Therefore, we use the Moore–Penrose inverse of $\tilde{X}_i^T X_i$ to obtain an estimate instead of (13), which is called "no regularization" in Fig. 4. The best regularization also scales as O(1/N) satisfying Theorem 1. Kernel-based regularization has the minimum MSE compared with other regularization forms because DC kernel utilizes the prior knowledge on the sparsity of coefficients. In addition, the MSEs of adaptive regularizations are always a little smaller than Tikhonov regularization and no regularization when $N < 10^5$.

Here we explain the reason why adaptive regularization with rank-1 kernel matrix fails to exhibit a clear advantage over typical non-adaptive protocol (as shown in Fig. 4) in the I.I. scenario. In the first step, for the chosen kernel matrix $S_i^{(1)}$, the condition $\theta_i \in R(S_i^{(1)}B)$ is usually not satisfied in the I.I. scenario. Thus, the estimate $\hat{\theta}_i^0$ is biased and MSE tends to a positive constant *c* as

$$\lim_{N \to \infty} \mathbb{E} \left\| \theta_i - \hat{\theta}_i^0 \right\| = c > 0.$$
(66)

Then in the second step, if we choose regularization as (35), $\theta_i \notin R(S_i^{\operatorname{rank-1}}B)$ because the only one vector in $R(S_i^{\operatorname{rank-1}}B)$ is $\hat{\theta}_i^0$ and $\lim_{N\to\infty} \mathbb{E} \left\| \theta_i - \hat{\theta}_i^0 \right\| = c > 0$. Moreover, even if we use multi-step regularization with rank-1 kernel matrix as above, the estimation result is still biased and MSE always tends to a constant, because the number of adaptive steps is always finite.



Fig. 4. The error scalings of different regularization forms with WLS using 10 types of 4 dimensional pure states. Except the best regularization, all the MSEs tend to constants as predicted by Theorem 1 and Remark 2 because $\theta_i \in R(S_iB)$ does not hold.

As *N* increases, except the best regularization, all the MSEs tend to constants as predicted by Theorem 1 because $\theta_i \in R(S_iB)$ does not hold.

5.2. Coherent states

Since coherent states are truncated, we consider a larger dimensional three-valued phase-sensitive detector as

$$P_1^{(8)} = U_1 \operatorname{diag} \left(P_1^{(4)}, P_1^{(4)} \right) U_1^{\dagger},$$

$$P_2^{(8)} = U_2 \operatorname{diag} \left(P_2^{(4)}, P_2^{(4)} \right) U_2^{\dagger},$$

$$P_3^{(8)} = I - P_1^{(8)} - P_2^{(8)},$$
(67)

where d = 8 and U_1 , U_2 are random unitary matrices (Johnston, 2016; Zyczkowski & Kus, 1994). We also ensure $P_3^{(8)}$ is positive semidefinite.

Since coherent states are more similar to each other, we generate 640 random different types of coherent states using the probe state preparation in Wang et al. (2021) where the real part and imaginary part of α are randomly generated in the interval [-1, 1]. We use different regularization forms including no regularization ((23) with c = 0), Tikhonov regularization ((23) with c = 10), kernel-based regularization ((28) with c = 0.2, $\mu = 0.9$), rank-1 adaptive regularization, full-rank adaptive regularization (see Section 4) and the best regularization (33). For rank-1 adaptive regularization, we use kernel-based regularization ((28) with c = 0.2, $\mu = 0.9$) in step 1 and (35) in step 2. For full-rank adaptive regularization, we use kernel-based regularization ((28) with c = 0.2, $\mu = 0.9$) in step 1 and (36) in step 2. For each resource number, we run the algorithm 100 times and obtain the average MSE and standard deviation.

The results are shown in Fig. 5. When $N < 10^8$, the MSEs of kernel-based regularization and adaptive regularization are a little smaller than Tikhonov regularization and no regularization. In addition, full-rank adaptive regularization has a little smaller MSE than rank-1 adaptive regularization. When $N > 10^8$, all the MSEs scale as O(1/N) satisfying Theorem 1. Since these coherent states are I.C., we also consider resource distribution optimization without regularization. The simulation results are shown in Fig. 6. We can also obtain a lower MSE with resource distribution optimization and both MSEs scale as O(1/N) for $N > 10^7$. Then using the same algorithm in Wang et al. (2021), we generate only 48 random types of coherent states where real parts and imaginary parts of α are randomly generated in the interval [-1, 1]. We use the same regularization (28) and the results are shown in Fig. 7. Kernel-based regularization and adaptive regularization always have smaller MSEs compared with Tikhonov regularization and



Fig. 5. The error scalings of different regularization forms with WLS using 640 types of coherent states. When $N > 10^8$, all the MSE scales as O(1/N) satisfying Theorem 1.



Fig. 6. The MSE comparison between average and optimized resource distribution using 640 types of coherent states.



Fig. 7. The error scalings of different regularization forms with WLS using 48 types of coherent states. Except the best regularization, all the MSEs tend to constants as predicted by Theorem 1 and Remark 2 because $\theta_i \in R(S_iB)$ does not hold.

no regularization. When $N > 10^{10}$, except the best regularization, all the MSEs tend to constants as predicted by Theorem 1 because $\theta_i \in R(S_iB)$ does not hold.

6. Experimental examples

We consider the same quantum optical experimental system for QDT in Yokoyama et al. (2019) and Wang et al. (2021). The entire experimental setup is given in Fig. 8. Wang et al. (2021) considered experiments for two different sets of detectors, denoted as Group I and Group II, respectively, and the basis of the POVM elements is the two-mode Fock state basis as $\{|0, 0\rangle, |1, 0\rangle, |0, 1\rangle, |2, 0\rangle, |1, 1\rangle, |0, 2\rangle$. The true values of Group



Fig. 8. The detailed structure and description for experimental setup can be found in Wang et al. (2021) and Yokoyama et al. (2019). Att., Attenuator; PBS, Polarization beam splitter; H, Half wave plate; Q, Quarter wave plate.



Fig. 9. Experimental and simulation QDT results of Tikhonov regularization (LS), rank-1 adaptive regularization (WLS) and full-rank adaptive regularization (WLS) for Group I.

I and II can be found in Wang et al. (2021). Wang et al. (2021) recorded 10^6 measurement outcomes for each input state, and repeated the process 6 times. We use these measurement data to identify the detectors and also plot the identification results using simulated measurement data as a comparison in Figs. 9 and 10.

For the QDT problem, Wang et al. (2021) employed Tikhonov regularization with standard LS estimation, where they chose $D_i^{\text{Tikhonov}} = \frac{10^3}{N}I$ and the estimation is given in (24), while here we use rank-1 adaptive regularization and full-rank adaptive regularization is similar to adaptive regularization, we only show the results of adaptive regularization.

To determine the hyper-parameters in the DI kernel, we use 15 probe states as estimation data and 4 probe states as validation data. In Group I, we choose c = 0.001, $\mu = 0.8$ in (28) in step 1 and (35) in step 2 for-rank 1 adaptive regularization and for full-rank adaptive regularization, we choose $c = 0.001, \mu =$ 0.8 in (28) in step 1 and (36) in step 2. The results are shown in Fig. 9. Adaptive regularization (WLS) performs better than Tikhonov regularization (LS) in Wang et al. (2021), especially for large resource number N. In addition, the MSE of full-rank adaptive regularization is a little smaller than rank-1 adaptive regularization. In Group II, for rank-1 adaptive regularization, we choose $c = 0.0008, \mu = 0.9$ in (28) in step 1 and (35) in step 2 and for full-rank adaptive regularization, we choose c =0.0008, $\mu = 0.9$ in (28) in step 1 and (36) in step 2. The results are shown in Fig. 10. Adaptive regularization (WLS) performs better than Tikhonov regularization (LS) when $N > 10^{2.5}$ and the MSE of full-rank adaptive regularization is always a little smaller than rank-1 adaptive regularization. Moreover, the MSE of Group II is a little larger than that of Group I because the amplitudes of nondiagonal elements in Group II are significantly larger than zero.



Fig. 10. Experimental and simulation QDT results of Tikhonov regularization (LS), rank-1 adaptive regularization (WLS) and full-rank adaptive regularization (WLS) for Group II.

7. Conclusion

In this paper, using regularization, we improve QDT accuracy with given probe states. In the I.C. and I.I. scenarios, we have employed WLS estimation, discussed different regularization forms, proved the scaling of MSE under the static assumption and characterized the best regularization. In the I.C. scenario, without regularization, we have studied resource distribution optimization and converted it to an SDP problem. The numerical examples have demonstrated the effectiveness of different regularization forms and resource distribution optimization. In a quantum optical experiment, our adaptive regularization with WLS has achieved lower mean squared errors compared with Tikhonov regularization with LS. It remains an open problem how to choose the kernel optimally in adaptive regularization.

Acknowledgment

S. Xiao would like to thank Dr. Xueke Zheng for the helpful discussion.

References

- Boyd, S., & Vandenberghe, L. (2004). Convex optimization. Cambridge University Press.
- Brida, G., Ciavarella, L., Degiovanni, I. P., Genovese, M., Lolli, L., Mingolla, M. G., et al. (2012). Quantum characterization of superconducting photon counters. *New Journal of Physics*, 14(8), Article 085001.
- Brida, G., Ciavarella, L., Degiovanni, I. P., Genovese, M., Migdall, A., Mingolla, M. G., et al. (2012). Ancilla-assisted calibration of a measuring apparatus. *Physical Review Letters*, 108, Article 253601.
- Burgarth, D., & Yuasa, K. (2012). Quantum system identification. *Physical Review Letters*, 108, Article 080502.
- Busch, P. (1991). Informationally complete sets of physical quantities. International Journal of Theoretical Physics, 30(9), 1217–1227.
- Cai, T., Kim, D., Wang, Y., Yuan, M., & Zhou, H. H. (2016). Optimal large-scale quantum state tomography with Pauli measurements. *The Annals of Statistics*, 44(2), 682–712.
- Chen, T. (2018). On kernel design for regularized LTI system identification. *Automatica*, 90, 109–122.
- Chen, T., & Andersen, M. S. (2021). On semiseparable kernels and efficient implementation for regularized system identification and function estimation. *Automatica*, 132, Article 109682.
- Chen, T., Andersen, M. S., Ljung, L., Chiuso, A., & Pillonetto, G. (2014). System identification via sparse multiple kernel-based regularization using sequential convex optimization techniques. *IEEE Transactions on Automatic Control*, 59(11), 2933–2945.

- Chen, T., & Ljung, L. (2013). Implementation of algorithms for tuning parameters in regularized least squares problems in system identification. *Automatica*, 49(7), 2213–2220.
- Chen, T., Ohlsson, H., & Ljung, L. (2012). On the estimation of transfer functions, regularizations and Gaussian processes—Revisited. *Automatica*, 48(8), 1525–1535.
- Cui, J., Li, C.-K., & Sze, N.-S. (2017). Products of positive semi-definite matrices. Linear Algebra and its Applications, 528, 17–24.
- D'Ariano, G. M., Maccone, L., & Presti, P. L. (2004). Quantum calibration of measurement instrumentation. *Physical Review Letters*, 93, Article 250407.
- Degen, C. L., Reinhard, F., & Cappellaro, P. (2017). Quantum sensing. Reviews of Modern Physics, 89, Article 035002.
- DiVincenzo, D. P. (1995). Quantum computation. Science, 270(5234), 255-261.
- Dong, D., & Petersen, I. R. (2022). Quantum estimation, control and learning: Opportunities and challenges. Annual Reviews in Control, 54, 243–251.
- Eldar, Y. (2006). Uniformly improving the Cramér-Rao bound and maximumlikelihood estimation. *IEEE Transactions on Signal Processing*, 54(8), 2943–2956.
- Feito, A., Lundeen, J. S., Coldenstrodt-Ronge, H., Eisert, J., Plenio, M. B., & Walmsley, I. A. (2009). Measuring measurement: theory and practice. *New Journal of Physics*, 11(9), Article 093038.
- Fiurášek, J. (2001). Maximum-likelihood estimation of quantum measurement. *Physical Review A*, 64, Article 024102.
- Fiurášek, J., & Hradil, Z. (2001). Maximum-likelihood estimation of quantum processes. *Physical Review A*, 63, Article 020101.
- Grandi, S., Zavatta, A., Bellini, M., & Paris, M. G. A. (2017). Experimental quantum tomography of a homodyne detector. *New Journal of Physics*, 19(5), Article 053015.
- Grant, M., & Boyd, S. (2008). Graph implementations for nonsmooth convex programs. In V. Blondel, S. Boyd, & H. Kimura (Eds.), *Lecture notes in control and information sciences*, *Recent advances in learning and control* (pp. 95–110). Springer-Verlag Limited, http://stanford.edu/~boyd/graph_dcp. html.
- Grant, M., & Boyd, S. (2014). CVX: Matlab software for disciplined convex programming, version 2.1. http://cvxr.com/cvx.
- Guță, M., & Yamamoto, N. (2016). System identification for passive linear quantum systems. IEEE Transactions on Automatic Control, 61(4), 921–936.
- Hou, Z., Zhong, H. S., Tian, Y., Dong, D., Qi, B., Li, L., et al. (2016). Full reconstruction of a 14-qubit state within four hours. *New Journal of Physics*, 18(8), Article 083036.
- Johnston, N. (2016). QETLAB: A MATLAB toolbox for quantum entanglement, version 0.9. http://dx.doi.org/10.5281/zenodo.44637.
- Lundeen, J. S., Feito, A., Coldenstrodt-Ronge, H., Pregnell, K. L., Silberhorn, C., Ralph, T. C., et al. (2009). Tomography of quantum detectors. *Nature Physics*, 5(1), 27–30.
- Mazzoleni, M., Chiuso, A., Scandella, M., Formentin, S., & Previdi, F. (2022). Kernel-based system identification with manifold regularization: A Bayesian perspective. *Automatica*, 142, Article 110419.
- Miszczak, J. A. (2012). Generating and using truly random quantum states in mathematica. *Computer Physics Communications*, 183(1), 118–124.
- Mu, B., & Chen, T. (2018). On input design for regularized LTI system identification: Power-constrained input. *Automatica*, *97*, 327–338.
- Mu, B., Chen, T., & Ljung, L. (2018). On asymptotic properties of hyperparameter estimators for kernel-based regularization methods. *Automatica*, 94, 381–395.
- Mu, B., Chen, T., & Ljung, L. (2021). On the asymptotic optimality of crossvalidation based hyper-parameter estimators for regularized least squares regression problems. math.ST, arXiv:2104.10471.
- Mu, B., Qi, H., Petersen, I. R., & Shi, G. (2020). Quantum tomography by regularized linear regressions. *Automatica*, 114, Article 108837.
- Natarajan, C. M., Zhang, L., Coldenstrodt-Ronge, H., Donati, G., Dorenbos, S. N., Zwiller, V., et al. (2013). Quantum detector tomography of a timemultiplexed superconducting nanowire single-photon detector at telecom wavelengths. *Optics Express*, 21(1), 893–902.
- Nielsen, M. A., & Chuang, I. L. (2010). Quantum computation and quantum information. Cambridge University Press.
- Nurdin, H. I., & Guţă, M. (2022). Parameter estimation and system identification for continuously-observed quantum systems. *Annual Reviews in Control*, 54, 295–304.
- Pillonetto, G., Chen, T., Chiuso, A., De Nicolao, G., & Ljung, L. (2022). Regularized system identification: Learning dynamic models from data. Springer Nature.
- Pillonetto, G., Chen, T., Chiuso, A., Nicolao, G. D., & Ljung, L. (2016). Regularized linear system identification using atomic, nuclear and kernel-based norms: The role of the stability constraint. *Automatica*, 69, 137–149.

- Pillonetto, G., Chiuso, A., & De Nicolao, G. (2011). Prediction error identification of linear systems: A nonparametric Gaussian regression approach. *Automatica*, 47(2), 291–305.
- Pillonetto, G., & De Nicolao, G. (2010). A new kernel-based approach for linear system identification. Automatica, 46(1), 81–93.
- Pillonetto, G., Dinuzzo, F., Chen, T., Nicolao, G. D., & Ljung, L. (2014). Kernel methods in system identification, machine learning and function estimation: A survey. *Automatica*, 50(3), 657–682.
- Prugovečki, E. (1977). Information-theoretical aspects of quantum measurement. International Journal of Theoretical Physics, 16(5), 321–331.
- Qi, B., Hou, Z., Li, L., Dong, D., Xiang, G.-Y., & Guo, G.-C. (2013). Quantum state tomography via linear regression estimation. *Scientific Reports*, 3, 3496.
- Renema, J. J., Frucci, G., Zhou, Z., Mattioli, F., Gaggero, A., Leoni, R., et al. (2012). Modified detector tomography technique applied to a superconducting multiphoton nanodetector. *Optics Express*, 20(3), 2806–2813.
- Sone, A., & Cappellaro, P. (2017a). Exact dimension estimation of interacting qubit systems assisted by a single quantum probe. *Physical Review A*, *96*, Article 062334.
- Sone, A., & Cappellaro, P. (2017b). Hamiltonian identifiability assisted by a single-probe measurement. *Physical Review A*, 95, Article 022335.
- Wang, Y., Dong, D., Qi, B., Zhang, J., Petersen, I. R., & Yonezawa, H. (2018). A quantum Hamiltonian identification algorithm: Computational complexity and error analysis. *IEEE Transactions on Automatic Control*, 63(5), 1388–1403.
- Wang, Y., Dong, D., Sone, A., Petersen, I. R., Yonezawa, H., & Cappellaro, P. (2020). Quantum Hamiltonian identifiability via a similarity transformation approach and beyond. *IEEE Transactions on Automatic Control*, 65(11), 4632–4647.
- Wang, Y., Dong, D., & Yonezawa, H. (2019). Tomography of binary quantum detectors. In 2019 IEEE 58th conference on decision and control (pp. 396–400).
- Wang, Y., Yin, Q., Dong, D., Qi, B., Petersen, I. R., Hou, Z., et al. (2019). Quantum gate identification: Error analysis, numerical results and optical experiment. *Automatica*, 101, 269–279.
- Wang, Y., Yokoyama, S., Dong, D., Petersen, I. R., Huntington, E. H., & Yonezawa, H. (2021). Two-stage estimation for quantum detector tomography: Error analysis, numerical and experimental results. *IEEE Transactions* on Information Theory, 67(4), 2293–2307.
- Wu, P. Y. (1988). Products of positive semidefinite matrices. *Linear Algebra and its Applications*, 111, 53–61.
- Xiao, S., Wang, Y., Dong, D., & Zhang, J. (2022). Optimal and two-step adaptive quantum detector tomography. *Automatica*, *141*, Article 110296.
- Xiao, S., Xue, S., Dong, D., & Zhang, J. (2021). Identification of time-varying decoherence rates for open quantum systems. *IEEE Transactions on Quantum Engineering*, 2, Article 2102212.
- Yokoyama, S., Pozza, N. D., Serikawa, T., Kuntz, K. B., Wheatley, T. A., Dong, D., et al. (2019). Characterization of entangling properties of quantum measurement via two-mode quantum detector tomography using coherent state probes. *Optics Express*, 27(23), 34416–34433.
- Zhang, L., Coldenstrodt-Ronge, H. B., Datta, A., Puentes, G., Lundeen, J. S., Jin, X.-M., et al. (2012). Mapping coherence in measurement via full quantum tomography of a hybrid optical detector. *Nature Photonics*, 6(6), 364.
- Zhang, L., Datta, A., Coldenstrodt-Ronge, H. B., Jin, X.-M., Eisert, J., Plenio, M. B., et al. (2012). Recursive quantum detector tomography. *New Journal of Physics*, 14(11), Article 115005.
- Zhang, J., & Sarovar, M. (2014). Quantum Hamiltonian identification from measurement time traces. *Physical Review Letters*, 113, Article 080401.
- Zhang, J., & Sarovar, M. (2015). Identification of open quantum systems from observable time traces. *Physical Review A*, 91(5), Article 052121.
- Zyczkowski, K., & Kus, M. (1994). Random unitary matrices. Journal of Physics A: Mathematical and General, 27(12), 4235–4245.



Shuixin Xiao received a B.E. degree in automation from Huazhong University of Science and Technology, Wuhan, China, in 2018. He is currently working toward the Ph.D. degree in a joint Ph.D. program with Shanghai Jiao Tong University, Shanghai, China, and the University of New South Wales, Canberra, ACT, Australia. His current research interests include quantum system identification, quantum computing and quantum information processing.



Yuanlong Wang received a B.E. degree in automation from Northeastern University, Shenyang, China, in 2011, an M.E. degree in control science and engineering from Zhejiang University, Hangzhou, China, in 2015, and a Ph.D. degree in engineering from the University of New South Wales, Canberra, ACT, Australia, in 2019. He was a Post-Doctoral Fellow with the Centre for Quantum Dynamics, Griffith University, Brisbane, QLD, Australia, from 2019 to 2022. He is now an Associate Professor in the Academy of Mathematics and Systems Science, Chinese Academy of Sciences, Beijing, China.

His research interests include quantum control theory and quantum parameter estimation theory.



Jun Zhang received the B.S. degree in Automatic Control from Shanghai Jiao Tong University, China, in 1993, and the M.S. and Ph.D. degrees in Electrical Engineering from the University of California at Berkeley, CA, USA in 1999 and 2003, respectively. Later he worked as a research specialist at Berkeley Chemistry department, and also as a staff engineer at several high tech firms in the Silicon Valley. He is currently a Research Fellow at the Joint Institute of UMich-SJTU in Shanghai Jiao Tong University. His current research interests include quantum control, process and motion control.



Daoyi Dong is currently an Associate Professor at the University of New South Wales, Canberra, Australia. He received a B.E. degree in automatic control and a Ph.D. degree in engineering from the University of Science and Technology of China, Hefei, China, in 2001 and 2006, respectively. He was with the Institute of Systems Science, Chinese Academy of Sciences and with Zhejiang University. He had visiting positions at Princeton University, NJ, USA, RIKEN, Wako-Shi, Japan, University of Duisburg–Essen, Germany and The University of Hong Kong, Hong Kong.

His research interests include quantum control and machine learning. Dr. Dong was awarded an ACA Temasek Young Educator Award by The Asian Control Association and is a recipient of a Future Fellowship, an International Collaboration Award and an Australian Post-Doctoral Fellowship from the Australian Research Council, and a Humboldt Research Fellowship from the Alexander von Humboldt Foundation of Germany. He was a Member-at-Large, Board of Governors, and the Associate Vice President for Conferences & Meetings, IEEE Systems, Man and Cybernetics Society. He served as an Associate Editor of IEEE Transactions on Neural Networks and Learning Systems (2015–2021). He is currently an Associate Editor of IEEE Transactions on Mechatronics. He is an IEEE Fellow.



Shota Yokoyama received a B.E. degree, a M.E. degree and a Ph.D. in Engineering at the Department of Applied Physics, the University of Tokyo, Japan, in 2010, 2012 and 2015. He was a JSPS Postdoctoral Fellow for Research Abroad at the Japan Society for the Promotion of Science from 2015 to 2017. He has joined the School of Engineering and Information Technology at the University of New South Wales, Canberra, Australia, as a Visiting Fellow (2015–2017) and as a Research Associate since 2017. His current research interests are in experimental quantum optics, quantum information

and quantum computation.



Ian R. Petersen was born in Victoria, Australia. He received a Ph.D. in Electrical Engineering in 1984 from the University of Rochester, USA. From 1983 to 1985 he was a Postdoctoral Fellow at the Australian National University. In 1985 he joined the University of New South Wales, Canberra, Australia. He moved to The Australian National University in 2017 where he is currently a Professor in the School of Engineering. He was the Australian Research Council Executive Director for Mathematics, Information and Communications in 2002 and 2003. He was Acting Deputy Vice-Chancellor

Research for the University of New South Wales in 2004 and 2005. He held an Australian Research Council Professorial Fellowship from 2005 to 2007, an Australian Research Council Federation Fellowship from 2007 to 2012, and an Australian Research Council Laureate Fellowship from 2012 to 2017. He served as an Associate Editor for the IEEE Transactions on Automatic Control, Systems and Control Letters, Automatica, and SIAM Journal on Control and Optimization. Currently he is an Editor for Automatica in the area of optimization in systems and control. He is a fellow of IFAC, the IEEE and the Australian Academy of Science.

His main research interests are in robust control theory, quantum control theory and stochastic control theory. Ian Petersen was elected IFAC Council Member for the 2014–2017 Triennium. He was also elected to be a member of the IEEE Control Systems Society Board of Governors for the periods 2011–2013 and 2015–2017. He was Vice-president for Technical Activity for the Asian Control Association and was General Chair of the 2012 Australia Control Conference. He was General Chair of the 2015 IEEE Multi-Conference on Systems and Control.



Hidehiro Yonezawa received his B.E. degree, M.E. degree and Ph.D. in Engineering at the Department of Applied Physics, the University of Tokyo, Japan, in 2002, 2004 and 2007. Upon completion of his Ph.D., he was appointed as Research Associate (2007–2009) and then promoted to Project Assistant Professor (2009–2013) in the University of Tokyo. He commenced his current appointment as a Senior Lecturer at the University of New South Wales Canberra in September 2013. He has joined the Centre for Quantum Computation and Communication technology (CQC2T) as a Co-Program

Manager since 2015. Through his career he has been working on experimental quantum optics, quantum information and quantum control.